

# **A new analytical model for the response of AlGaN/GaN HEMT‑based pH sensors**

**Kavita Thorat Upadhyay1,2 · Manju K. Chattopadhyay3**

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#### **Abstract**

pH sensors are monitoring devices with wide applications in biology, chemistry, medicine, and agriculture. To enhance their sensitivity and long-term stability, efficient study of such devices becomes imperative but is impossible without the aid of accurate analytical models. A new analytical model for the pH sensing characteristics of AlGaN/GaN-based high-electronmobility transistors (HEMTs) is presented herein, as well as theoretical predictions and optimization of the charge sensitivity for ungated AlGaN/GaN HEMT-based pH sensors. The change in the drain current with the changing surface potential due to a variation in the pH of the electrolyte is calculated for devices with diferent Al mole fractions, AlGaN thicknesses, gate length spacings, and passivation layers. The numerical values for the drain current, threshold voltage, and surface potential obtained by using this new model show good agreement with available experimental results. It is demonstrated that the sensitivity of GaN HEMT-based pH sensors at lower pH values can be improved by applying a  $\text{SiN}_x$  passivation layer to the HEMT. The calculated average root-mean-square error of our model is 0.018, being an order of magnitude lower than other models reported in literature.

**Keywords** GaN HEMT · 2DEG · Analytical model · Chemical sensor · pH sensing · Sensitivity

# **1 Introduction**

The normal range of pH variation for human blood is 7.35–7.45. Any deviation in the pH value of a biomolecule may indicate a disease, possibly with severe effects [1]. pH measurement is also vital in many other felds including biomedicine, food and nutrition, ecological studies, chemistry, and oceanography. Electrochemical (EC) liquid sensors thus fnd application in numerous biomedical practices [2]. The absence of a gate electrode makes ion-sensitive feld-efect transistor (ISFET)-based sensors an important type of EC sensor for measuring pH  $[3, 4]$ . However, since the gate region in an ISFET is prone to charging efects, a reference electrode must be used to bias the gate area to above the threshold voltage and thus allow a drain current fow. Highelectron-mobility transistors (HEMTs), on the other hand, do not need a gate voltage to turn on.

A suitable choice of sensing material is as important as the device selection. During the last three decades, substantial efforts have been made towards the development of novel semiconductor materials exhibiting physical and mechanical properties that can overcome the fundamental limits of silicon-based conventional electronics, which are mature and less expensive. Among these materials, wide-bandgap (WBG) semiconductors such as silicon carbide (SiC)  $[5, 6]$ , diamond [7–9], and gallium nitride (GaN) have been shown to possess superior properties [10]. SiC and GaN are the most well-known WBG devices, offering attractive features that make them suitable for operation in high-temperature conditions, such as a wide bandgap (3 eV), high drift saturation velocity, high thermal conductivity, low intrinsic carrier concentration, and large critical electric feld. Both GaN and SiC belong to the same family of WBG semiconductors and share similar, attractive properties. SiC has attracted great attention for use in high-temperature applications [11].

 $\boxtimes$  Manju K. Chattopadhyay mkorwal@yahoo.com

<sup>1</sup> Department of Electronics and Communication Engineering, IPS Academy, Institute of Engineering and Science, Indore 452012, India

<sup>&</sup>lt;sup>2</sup> Department of Electronics and Telecommunication Engineering, Institute of Engineering and Technology, Devi Ahilya University, Indore 452017, India

<sup>&</sup>lt;sup>3</sup> School of Electronics, Devi Ahilya University, Khandwa Road, Indore 452017, India

However, most work reported in this feld has been realized with either large areas or low device density integration. The direct temperature dependence of the carrier concentration, due to the bulk nature of the active region in SiC devices, is a known shortcoming of SiC [12]. Moreover, at higher temperatures, the junction leakage of SiC devices is degraded due to crystal dislocations. GaN offers a considerable performance improvement over SiC with respect to the response speed, operating temperature limits, and reverse recovery characteristics [12]. At low breakdown voltages, GaN devices offer lower channel resistance than SiC devices. Owing to its comparatively better physical and chemical stability, GaN is also an ideal material for the development of explicit chemical sensor and biosensor systems [13]. A detailed review on signifcant research work in the feld of sensor technology based on GaN heterostructures can be found in our recent work [14].

AlGaN/GaN HEMTs have become popular for use in high-frequency radiofrequency (RF) power amplifer and power switching applications as well [15]. Their high sheet carrier concentration derives from the polarization-induced two-dimensional electron gas (2DEG) produced at the AlGaN–GaN interface. The density of the 2DEG is balanced by the surface charge state. The sheet charge density in the 2DEG channel is thus very sensitive to the charge on the AlGaN surface. The strong chemical stability and the WBG of group III N materials ensure that such AlGaN/GaN heterostructures can operate as chemical sensors in harsh acidic or alkaline environments and at very high temperatures. In particular, these chemical sensors can be incorporated into biosystems because of their biological compatibility and nontoxicity. This has provided the foundation for new acute medical care systems in recent years [16, 17].

As stated above, pH measurement is a major consideration in almost all felds involving chemicals. pH monitoring is also required for seawater analysis, for soil analysis to check its fertility, and in the food processing industry. pH also plays a decisive role in targeted drug release from nanoscale drug carriers. Based on these applications, pH measurement has become an extremely important topic for researchers working in diverse felds [18].

#### **1.1 The sensing mechanism**

An electrolyte solution with diferent pH values can be considered to represent a type of intrinsic semiconductor material because the mobile ions in solution are similar to electron and hole charges in semiconductors. The pH sensing mechanism is related to reactions between ions in the electrolyte solution with diferent pH values and the positive surface charge induced by the polarization on the surface of HEMTs. A probable mechanism for the adsorption of positive and negative charges that change the potential at the AlGaN/Au and Au/electrolyte interfaces related to the  $H<sup>+</sup>$  concentration is explained by the site binding model introduced by Yates et al. and Munch et al. [19, 20]. In this model, hydroxyl groups at the surface act as amphoters. They may be dependent on the  $H<sup>+</sup>$  concentration and the equilibrium constants for the related reactions:

$$
MOH \leftrightarrow MO^{-} + H^{+}
$$
 (1)

$$
MOH + H^{+} \leftrightarrow MOH_{2}^{+}
$$
 (2)

Here, M is either Si or a metal in MOH. Because of the formation of more protonized hydroxyls  $MOH<sub>2</sub><sup>+</sup>$ , the 2DEG concentration increases to balance the positive charge induced on the surface. Vice versa, negative charges at the insulator surface decrease the  $H<sup>+</sup>$  concentration in solution. A high sensitivity of GaN-based pH sensors to hydronium ions was reported in Ref. [21]. The sensitivity *S* of such sensors can be regarded as capturing the modifcation of the surface potential  $\psi_0$  as a function of the concentration of hydronium ions in the electrolyte.

In diferent electrochemical and non-electrochemical methods used for measuring pH, AlGaN/GaN HEMTs have shown great promise for achieving unprecedented speed and sensitivity. Indeed, GaN-based pH sensors showed high sensitivity of 57.3 mV/pH to hydronium ions [22, 23]. The sensitivity of these sensors is related to the action between ions in the electrolyte on the open region of the device and the positive surface charge induced by polarization. This action causes a change in the surface charge on the devices. Nonmodulation-doped GaN HEMTs have been investigated due to their advantages such as a reduced gate leakage current, lower  $V_{\text{pinch-off}}$  value, and lower noise due to the absence of donors from the AlGaN layer [24, 25].

Stutzmann et al. [26] and Mehandru et al. [27] reported preliminary results on the response of open-gate AlGaN/ GaN HEMTs to polar liquids. Kokowa et al. first experimentally investigated the pH-sensing characteristics of unpassivated u-AlGaN/GaN HEMT structures [28]. Later, Abidin et al. [24] investigated the surface of the open gate of an HEMT experimentally. An increase in the pH level leads to a shift of the threshold voltage towards the positive side.

Sharma et al. verifed the characteristics of a pH and salinity sensor fabricated using gated AlGaN/GaN HEMT structures in phosphate buffer saline (PBS) and aqueous salt solutions [29]. In deionized (DI) water, the HEMT devices displayed  $I_d$ – $V_d$  characteristics similar to the output characteristics of classical HEMT structures in air. A sensitivity of 6.48 mA/mm-molar and a response time of 250–350 ms at  $V_{ds}$  = +1 V were obtained using these GaN HEMT structures [29]. Cheng et al. demonstrated a pH sensor based on a planar dual-gate AlGaN/GaN HEMT cascode amplifer that enhanced the pH sensitivity by about 45 times from 45 mV/pH to 2.06 V/pH with linearity of 1.27%  $[30]$ . Xue et al. illustrated how to adjust the threshold voltage  $V<sub>T</sub>$  of an AlGaN/GaN HEMT-based pH sensor from −3.46 to −1.15 V by applying photo-electrochemical (PEC) oxidation on the GaN cap layer surface. They postulated that the sensitivity of this reference electrode sensor could be considerably improved by varying the  $V_T$  value so that  $V_G|g_{mMAX}$ would approach the corresponding gate voltage when a droplet is placed on the sensing window plane. This approach may be benefcial for scaling down and integration of future AlGaN/GaN HEMT pH sensors [31].

Zhang et al. theoretically and experimentally investigated the result of varying the gate geometry of a pH sensor on its sensitivity. The series resistance  $(R<sub>S</sub>)$  of the packaged sensor was found to be an important feature limiting the current sensitivity. They found an optimum for  $W/L = \rho_{(2DEG)}/R_S$ [32].

#### **1.2 The analytical model**

Since the formation of the 2DEG determines the drain current fowing in a HEMT device, an accurate physics-based analytical expression for the 2DEG electron sheet concentration  $(n_s)$  is a prime necessity for developing a model for such devices. The main complexity in analytical modeling of the 2DEG electron sheet concentration  $(n<sub>s</sub>)$  is its subtle variation with the applied voltage.

Several theoretical–mathematical models for the drain *I*–*V* characteristics of AlGaN/GaN HEMT devices have been reported in literature [33–40], among which the model described by Toufk et al. [34] is one of the most popular. This is an empirical method that uses a large number of ftting parameters that must be extracted from experimental data. However, physics-based models are preferred for better prediction of the statistical variations in the device [35]. They follow well-defned geometrical and temperature scaling rules [36–39]. Khandelwal et al. [41] proposed a model for  $n_s$  based on only the first quantum energy level  $E_0$  of the conduction band of GaN. That model does not account for the electron concentration in the AlGaN layer and uses an interpolation function for  $n_s$ . Karumuri et al.  $[42]$  divided the operating zone of the HEMT into three regions depending on the relative position of the Fermi level  $E_F$ : region 1 for  $E_F$  <  $E_0$ , region 2 for  $E_0$  <  $E_F$  <  $E_C$ , and region 3 for  $E_F$  >  $E_C$ . In region 3, the  $E_F$  crosses the conduction-band minimum of AlGaN at the interface. The electron concentration in the AlGaN barrier layer  $(n<sub>B</sub>)$  thus also becomes important for the accuracy of the 2DEG model. Both Karumuri et al. [43] and Swamy et al. [43] developed models using simplifed Fermi–Dirac statistics. Swamy et al. [43] proposed a model for  $n<sub>s</sub>$  that does not need an interpolation function and uses only four ftting parameters. In contrast, the model developed by Khandelwal et al. [41] uses two interpolation functions.

The accuracy level of the model proposed by Swamy et al. [43] is comparable to that of Karumuri et al. [42].

Rabbaa et al. [44] applied a basic theoretical model to calculate the 2DEG electron sheet concentration  $(n_s)$  and the drain current in the HEMT as well as the surface potential due to the pH of the electrolyte solution. However, their model uses iterative techniques that are unsuitable for circuit simulations. Their method of calculating the saturation velocity requires excessive computation, too. Their model suffers from the additional limitation that the parameters used must be fne-tuned by approximately 5% to obtain a perfect match with experimental results.

To overcome the above-mentioned problems, we modified the model for  $n_s$  proposed by Swamy et al. [43] for quaternary alloy AlInGaN/GaN HEMTs and obtained a satisfactory match of the results with experimental data [45]. In the present work, we extend this modifed model to make it applicable for the investigation of pH sensing applications using AlGaN/GaN HEMTs, making the following changes:

- (i) The electron concentration  $n<sub>B</sub>$  in the AlGaN barrier layer is also included in the unifed model of Ref. [43]. Here, since we are analyzing a model for pH sensors operating in the region with  $E_F < E_c$  and with a negligible  $n<sub>B</sub>$ , we use the model only for  $n<sub>s</sub>$  but deliberately neglect  $n<sub>B</sub>$ , resulting in a simplified model. Thus, our parameters are chosen diferent from those in Swamy's model.
- (ii) Iterative techniques are not used to calculate the surface potential caused by the pH of the electrolyte solution.
- (iii) The values of the parameters are not approximated as done by Rabbaa et al. [44]. We use suitable ftting parameters without approximation.
- (iv) The model of the drain current is not divided into regions, unlike that of Rabbaa et al. [44], thus simplifying it further.
- (v) The current model is implemented specifically for AlGaN/GaN pH sensors.

The results calculated using the new model show good agreement with experimental results reported in literature. This paper is organized as follows: Sect. 2 discusses the device structure. Sections 3.1, 3.2, and 3.3 present the basis for our model for  $n_s$  and the charge-based current models. Section  $3.4$ introduces the model for the surface potential. Section 4 presents the results, while conclusions are drawn in Sect. 5.

## **2 The device structure**

A cross-sectional view of the AlGaN/GaN HEMT device structure is shown in Fig. 1. This structure consists of a thick GaN buffer layer with a width of about  $1-2 \mu m$  and

an AlGaN barrier layer with a typical thickness ranging between 10 and 30 nm. The gate bias can be applied at the open-gate area through a saturated calomel electrode (SCE), which is the reference electrode having a well-defned stable potential that is not afected by the concentration of hydrogen ions in the sample solution. To measure the pH value, the device along with the SCE is placed into an electrochemical cell flled with electrolyte solution [28].

Conducting channels form in the HEMTs due to the conducting layer of free electrons at the heterointerface between the doped wide-bandgap semiconductor (AlGaN) and the undoped narrower-bandgap semiconductor (GaN). Electrons from the edge region of the doped AlGaN are transferred into the lower-bandgap GaN semiconductor and collect in the quantum well that forms in the GaN at the junction of the two materials. The electrons are confned in the quantum well and are known as a 2DEG. The 2DEG is located in the region of undoped GaN, where the carrier mobility is high due to the reduced scattering efect. The 2DEG offers the advantage that there are no impurity atoms in the undoped GaN quantum well. Changes in the heterojunction space result in changes in the 2DEG concentration, thereby afecting the source–drain current. Due to the high carrier concentration, high electron mobility, and wide bandgap, such HEMT devices are sensitive to noise and even small variations in the current magnitudes.

# **3 A description of the model**

It has been established that the formation of a 2DEG in the heterostructure is the most important effect of the polarization property of nitride alloys. Literature presents many useful predictions based on theoretical and experimental studies of the polarization property of ternary AlGaN alloys. For the AlGaN/GaN HEMT, the threshold voltage  $V_{th}$  is related to the polarization as [44],



**Fig. 1** A cross-sectional view of the open-gate AlGaN/GaN HEMT dipped in an electrolyte. The gate bias is applied through the gate electrode to the electrolyte–AlGaN interface

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$$
V_{\text{th}} = \frac{\varnothing_{\text{b}}}{q} - \frac{\Delta E_{\text{c}}}{q} - \frac{qn_{\text{d}}d^2}{2\varepsilon} - \frac{q\sigma_{\text{int}}d}{\varepsilon} \tag{3}
$$

where  $\varnothing_b$  is the height of the Schottky barrier,  $\Delta E_c$  is the conduction-band discontinuity,  $q$  is the electron charge,  $n_d$  is the doping concentration,  $\varepsilon$  is the permittivity of the material, *d* is the thickness of the AlGaN layer,  $\sigma_{int} = P_{Gal} - P_{AIGaN}$ where  $P_{\text{GaN}}$  is the polarization of GaN and  $P_{\text{AlGaN}}$  is the polarization of AlGaN, given by the sum of the spontaneous and piezoelectric polarizations:

$$
P_{\text{AlGaN}} = P_{\text{sp}}(\text{AlGaN}) + P_{\text{pz}}(\text{AlGaN}) \tag{4}
$$

## **3.1 The carrier density model**

The sheet charge density  $n_s$  of the 2DEG as a function of position  $x$  in the channel can be obtained by solving the onedimensional Poisson equation and is given as [44]:

$$
n_{\rm s} = \frac{\varepsilon}{qd_{\rm i}} \left( V_{\rm g} - V_{\rm th} - V_{\rm y} - \frac{E_{\rm F}}{q} \right),\tag{5}
$$

where  $d_i$  is the total length of the AlGaN and GaN layer,  $V_g$ is the gate voltage,  $V_y$  is the channel potential at a distance *y*, and  $E_F$  is the Fermi level.

As mentioned above, the current model is used to analyze an AlGaN/GaN heterostructure operating as a pH sensor in the region  $E_F$ < $E_c$ , where the electron concentration in the AlGaN layer  $n<sub>B</sub>$  is negligible. Hence, only  $n<sub>s</sub>$  is used in the model while  $n<sub>B</sub>$  is neglected.  $n<sub>s</sub>$  can be expressed analytically as a function of  $V_g$  thus [45]

$$
n_{s} = \frac{AV_{g0}}{(1+B)} \left( \frac{1 - A^{\frac{2}{3}} \gamma_{0}}{(1+B)^{\frac{2}{3}} V_{g0}^{\frac{1}{3}}} \right),
$$
 (6)

where  $V_{\rm go} = V_{\rm g} - V_{\rm th}$ ,  $A = \varepsilon / qd$ ,  $B = A/qD$ ,  $D$  is the density of states, and *γ*0 is a constant estimated from Shubnikov–De Haas or cyclotron resonance experiments.

## **3.2 The charge‑based drain current model**

An analytical model for the current is formulated using the defnition of the drain current along the channel given by

$$
I_{\rm D} = q\mu W \left[ n_T \left( \frac{\mathrm{d}V_y}{\mathrm{d}y} \right) - \left( \frac{kT}{q} \right) \left( \frac{\mathrm{d}n_T}{\mathrm{d}y} \right) \right],\tag{7}
$$

$$
dV_{y} = -dn_{T}\bigg(\frac{A+D}{AD} + \frac{2}{3}\gamma_{0}n_{T}^{\frac{-1}{3}}\bigg),\tag{8}
$$

where *W* is the width,  $\mu$  is the low-field mobility of the device,  $n<sub>T</sub>$  is the sheet carrier density in the channel that

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contributes to current conduction at any given distance *y*, *k* is the Boltzmann constant, and *T* is the ambient temperature. Taking the limit of *y* from the source to drain and integrating Eq. 7 yields the following expression for the drain current:

$$
I_{\rm D} = \frac{q\mu_{\rm eff}W}{L_{\rm g}} \left[ \frac{A+D}{2AD} \left( n_{\rm s}^2 - n_{\rm D}^2 \right) + \frac{2}{5} \gamma_0 \left( n_{\rm s}^{\frac{5}{3}} - n_{\rm D}^{\frac{5}{3}} \right) + \frac{kT}{q} \left( n_{\rm s} - n_{\rm D} \right) \right],\tag{9}
$$

where  $n_s$  and  $n_D$  are the charge carrier concentration at the source and the drain, respectively.  $\mu_{\text{eff}}$  is the effective 2DEG mobility, and  $L<sub>g</sub>$  is the gate length.  $n<sub>s</sub>$  here is calculated using Eq. (6). The calculation of  $n<sub>D</sub>$  is explained in the next section.

#### **3.3 The saturation voltage**

The saturation voltage is modeled using the approach given in Ref. [46]:

$$
V_{\rm sat} = \frac{v_{\rm s} V_{\rm go}}{v_{\rm s} + \frac{\mu_{\rm eff} V_{\rm go}}{2L}},\tag{10}
$$

where  $v<sub>s</sub>$  is the saturation velocity. The above Eq. (10), commonly used in many metal–oxide–semiconductor feld-efect transistor (MOSFET) models, is adopted here for the case of HEMTs to obtain the efective mobility as [46]

$$
\mu_{\text{eff}} = \frac{\mu_0}{\left(1 + a_1 V_{\text{go}} + a_2 V_{\text{go}}^2\right) \left(1 + a_3 V_{\text{ds}}\right)},\tag{11}
$$

where  $a_1$ ,  $a_2$ , and  $a_3$  are fitting parameters.

Once the saturation voltage is calculated, the effective drain voltage  $V_{\text{eff},D}$  can be obtained.  $V_{\text{eff},D}$  is designed to achieve a smooth transition between the applied drain–source voltage  $V_{ds}$  and the saturation voltage  $V_{sat}$  [43]:

$$
V_{\text{eff,D}} = V_{\text{sat}} \left[ 1 - \frac{\ln \left[ 1 + exp(1 - \alpha \frac{V_{\text{ds}}'}{V_{\text{sat}}}) \right]}{\ln(1 + \epsilon)} \right],
$$
 (12)

where  $\alpha$  is the transition width parameter and  $V'_{ds}$  represents the effective drain–source voltage under the calculated by subtracting the voltage drops due to the resistances at the source and drain from the drain–source voltage. The charge carrier concentration at the drain  $n<sub>D</sub>$  is calculated using Eq. (6) by replacing  $V_{\text{go}}$  with  $V_{\text{gdo}} = V_{\text{go}} - V_{\text{eff,D}}$ . As mentioned above, for pH sensing applications, only the sheet charge carrier concentration  $n<sub>s</sub>$  need be considered but not  $n<sub>B</sub>$ . Hence, the values of the fitting parameters in our model are diferent from those of Ref. [43]. This makes our model simpler, too. The values of these parameters are taken as presented in Table 1, calculated for three diferent devices to compare the results presented herein with various experimental data and thereby validate the model. All the parameters for device 1 are the same as in Ref. [43] except for the transition width parameter *α*.

#### **3.4 pH sensing**

HEMTs show significant sensitivity to relatively small changes in the pH level of a liquid. When the surface of a HEMT is exposed to an electrolyte solution with a specifc concentration, the change in the surface charge and potential causes a change in the channel charge density, which further modulates the drain current. This mechanism can be used to build a sensor to measure the pH of the solution.

Therefore, the efect of the electrolyte solution on the device manifests in the potential [47], as follows:

$$
\psi_0 = -\log_e 10 \delta_t V_{th} (\text{pH} - \text{pH}_{PZC}),\tag{13}
$$

where

 $\delta_t = \frac{\gamma}{1+\gamma}, \gamma = \frac{qN_{ss}\eta}{C_{eq}V_{th}}, \eta = 2 \times 10^{-\frac{(pK_b - pK_a)}{2}}, C_{eq} = \frac{C_{DL}C_{Stern}}{C_{DL} + C_{Stern}},$  $C_{\text{DL}} = \frac{\sqrt{8\epsilon V_T q n_0}}{2V}$  $\frac{2V_T q n_0}{2V_{\text{th}}}$ ,  $pH_{PZC} = \frac{pK_a + pK_b}{2}$ . Here *γ* is defined by the Gouy–Chapman Stern model [47],  $C_{Stern}$  is the capacitance of the Stern layer  $(C_{\text{Stern}} = 20 \,\mu\text{F/cm}^2)$  [47],  $n_0$  is the ionic charge concentration in the electrolyte,  $V_T$  is the thermal voltage, and  $N_{ss}$  is the site binding charge, given in Table 2.

The surface potential  $\psi_0$  is subtracted from the threshold voltage to obtain the effective threshold voltage  $V_{\text{eth}}$ , represented as [44]

$$
V_{\text{eth}} = V_{\text{th}} - \Psi_0. \tag{14}
$$

This effective threshold  $V_{\text{eth}}$  replaces the threshold voltage  $V_{\text{th}}$  in all the calculations.

## **4 Results**

To validate the model, the results are compared with the experimental data for device 1 reported in Ref. [48]. Device 1 has a gate length of 1 μm and a gate width of 75 μm. The Al mole fraction is 15% in the 25-nm AlGaN layer of device 1. Figure 2 depicts the drain current versus the gate voltage to compare the values calculated using the presented model, those calculated using the model introduced in Ref. [43], and the experimental data from Ref. [48]. The results show that all three sets of results match very well for  $V_d = 5$  V, even though our model has fewer parameters than that described in Ref.  $[43]$  and  $n<sub>B</sub>$  is excluded as well.

The surface potential is calculated for two surfaces of the HEMT. The first is the bare AlGaN surface, while the second one is an SiN*x*-covered AlGaN surface. An unintentional thin oxide layer forms on the exposed AlGaN surface when

Parameter (units)	Given in Ref. $[43]$ for	The model parameters			
	device 1	Device 1	Device 2	Device 3	
Depicted in	[43]	Figure 2	Figures 5 and 6	Figures 7 and 8	
Al mole fraction in AlGaN (dimensionless)	0.15	0.15	0.23	0.25	
AlGaN layer thickness (nm)	25	25	22	25	
Gate length $(\mu m)$			10	40	
Gate width $(\mu m)$	75	75	500	490	
$V_{\text{off}}(V)$	$-2.9$	$-2.9$	$-4.9$	$-6.1$	
Low-field mobility, $\mu_0$ (cm <sup>2</sup> /V-s)	600	600	950	1860	
Fitting parameter, $a_1$ (1/V)	$5.5 \times 10^{-3}$	$5.5 \times 10^{-3}$	10	$60 \times 10^{-3}$	
Fitting parameter, $a_2$ (1/V <sup>2</sup> )	$1.5 \times 10^{-3}$	$1.5 \times 10^{-3}$	0.05	$15 \times 10^{-3}$	
Fitting parameter, $a_3$ (1/V)	$2 \times 10^{-3}$	$2 \times 10^{-3}$	0.34	0.001	
Source–gate region resistance, $R_s(\Omega)$	0.72	0.72	0.45	0.75	
Drain–gate region resistance, $R_d(\Omega)$	0.72	0.72	0.45	0.75	
Transition width parameter, $\alpha$ (dimensionless)	0.01	0.32	0.82	1	

**Table 1** The parameter values used in the model for ftting and for various aspects of the device design

**Table 2** The site-binding parameters of diferent passivation layers [47]

Passivation layer		$\mathbf{p}K_{\alpha}$	$pK_h$	$N_{\rm cc}$
$SiN_r$	7.5	$-2$		$3 \times 10^{14}$ cm <sup>-2</sup>
$Ga_2O_3$	10.2	4.81	10.91	$3 \times 10^{14}$ cm <sup>-2</sup>



**Fig. 2** The drain current  $I_d$  versus the gate voltage  $V_g$  for a drain voltage of  $V_d = 5$  V, enabling a comparison of the results calculated herein with the experimental data in Ref. [48] for device 1 and the theoretical results obtained using the model described in Ref. [43]. The inset shows the same data plotted on a logarithmic scale. Solid dots represent the experimental data, while hollow symbols with dotted lines show the data calculated using the model described in Ref. [43], and continuous lines represent the calculations using the model described herein

it comes into contact with air [44]. This unintentional  $Ga<sub>2</sub>O<sub>3</sub>$ layer changes the surface potential. This efect is calculated by considering the parameters related to  $Ga_2O_3$  [44].

Figure 3 shows the variation of the surface potential with the change in the pH value for  $Ga_2O_3$  and the SiN<sub>x</sub> layer on the AlGaN surface of the HEMT. The results calculated for  $Ga<sub>2</sub>O<sub>3</sub>$  and  $SiN<sub>x</sub>$  match with the experimental results given by Rabbaa et al. [44] and Kokawa et al. [28], respectively. The surface potential decreases linearly with the pH value from 2 to 12.

A change in the surface potential causes a shift in the threshold voltage, which affects the sensitivity of the device. The device sensitivity is defned as



**Fig. 3** The variation of the surface potential with the pH value, illustrating the agreement of the calculations using the present model, for  $SiN<sub>x</sub>$  and  $Ga<sub>2</sub>O<sub>3</sub>$  passivation layers, with the experimental results presented in Ref. [28] and [44], respectively

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sensitivity = 
$$
\frac{V_{\text{thpH1}} - V_{\text{thpH2}}}{p\text{H1} - p\text{H2}} = \frac{\Delta V_{\text{th}}}{\Delta p\text{H}}.
$$
 (15)

Figure 4 shows the variation in the sensitivity of the device with different passivation layers, viz. SiN<sub>x</sub> and  $Ga<sub>2</sub>O<sub>3</sub>$ . The results of the present calculations show that, for  $\sin N_x$  and  $Ga_2O_3$  layers, the maximum sensitivity is 0.062 V/ pH and 0.060 V/pH while the average sensitivity is 0.595 V/ pH and 0.058 V/pH, respectively. The sensitivity is approximately equal to the Nernstian response, i.e., 0.59 V/pH. Kokawa et al. [28] reported an experimental pH sensitivity of 0.0567 V/pH with an SiN*x* layer. We also infer here that the change in the threshold voltage of the device is higher for  $pH < 7$ , i.e., for acidic solutions, while the sensitivity is lower for basic solutions. We observe that the device with an SiN*x* passivation layer is more sensitive than the device with a  $Ga<sub>2</sub>O<sub>3</sub>$  layer for acidic solutions.

Figures 5 and 6 depict the  $I_d - V_g$  characteristics of HEMT device 2 when covered with electrolyte solution of diferent pH values. Device 2 is an SiN*x*-layered HEMT with an Al mole fraction of 23%, a 22-nm-thick AlGaN layer, a gate length of 10  $\mu$ m, and a gate width of 500  $\mu$ m. Figure 5 compares the results with the experimental results of Kokawa et al.  $[28]$  and the theoretical results obtained using the  $I_d$ model proposed by Rabbaa et al. [44]. The values calculated herein at  $V_{ds}=0.2$  V for different pH values are in good agreement with the experimental results of Kokawa et al. [28]. It can be observed that the results of our model provide a better match than those of Ref. [44] without requiring any approximation in the parameters. The curves represent the results calculated using our model, while the symbols in the graph denote the experimental results of Kokawa et al. The calculated results match well with the experimental data



**Fig. 4** The variation of the sensitivity with the pH value for devices with  $\sin X_x$  and  $Ga_2O_3$  passivation layers. The sensitivity is nearly equal to the Nernstian response of 0.59 V/pH



**Fig. 5** The  $I_d - V_g$  characteristics with  $V_{ds} = 0.2$  V at pH 4 for device 2, showing a comparison of the data calculated using our model or that described in Ref. [44] with the experimental data in Ref. [28]. The continuous lines represent the data calculated using our model, the dashed line represents the theoretical calculations of Ref. [44], while the solid triangles represent the experimental data from Ref. [28]



**Fig. 6** The  $I_d-V_g$  characteristics with  $V_{ds}=0.2$  V at pH 10 for device 2, showing the comparison of the data calculated using our model with the experimental data from Ref. [28]. The lines represent our calculated values, while the solid dots denote the experimental data

for gate voltages above  $-1.5$  V. The minor variations seen below this voltage level may be due to undefned parameters such as the source and drain resistances. Also, some of this mismatch can be attributed to the Schottky barrier height at the metal–AlGaN interface, which may difer depending on the type and quality of the ohmic contact formed by various processes [49]. The units of current in the plots of the drain–current characteristics are diferent in Figs. 5 and 6 to

enable a comparison of the results calculated herein with the experimental data available in literature.

Figures 7 and 8 show the  $I_d$ – $V_d$  characteristics of device 3, which is a bare AlGaN HEMT with an electrolyte solution at pH 1.7 or 11.9. The gate voltages are kept at 0 V and −2 V for both curves. As discussed above, bare AlGaN is covered with  $Ga<sub>2</sub>O<sub>3</sub>$  when exposed to air. The calculations therefore adopt the parameters for  $Ga_2O_3$ . This HEMT has a 25-nm-thick AlGaN layer with an Al content of 25%. The gate width and gate length are fxed at 490 and 40 μm, respectively. The good match with the experiment results of



**Fig. 7** The drain current  $I_d$  and drain voltage  $V_d$  characteristics of summarize them in Tables 3, 4, and 5. HEMT device 3 at pH 1.7, showing a comparison of the data calculated using the current model versus the experimental data from Ref. [24] at gate voltages of 0 V and −2 V



**Fig. 8** The drain current  $I_d$  and drain voltage  $V_d$  characteristics of HEMT device 3 at pH 11.9, showing a comparison of the data calculated using the current model with the experimental data from Ref. [24] at gate voltages of 0 V and −2 V

Abidin et al. [24] seen in Figs. 7 and 8 validates the model presented herein. It is also observed that the surface potential decreases with increasing pH value of the electrolyte. This causes a decrease in the drain current of the device. The slight mismatch observed at higher drain voltages between our calculated results and the experimental data in Figs. 7 and 8 may be due to self-heating efects, which are not incorporated herein.

Figure 9 depicts the value of the effective threshold voltage calculated using Eq. (14) for devices 2 and 3 with electrolyte solutions of diferent pH values. The efective threshold is calculated by subtracting the surface potential  $\psi_0$  (Eq. 13) from the threshold voltage (Eq. 3). As observed in Fig. 3, the surface potential decreases with increasing pH value, resulting in a shift of the efective threshold value towards positive values.

The evaluation of the best-performing model relies on measurement error analysis [50]. The root-mean-square error (RMSE) is the most widely reported measurement error analysis parameter, representing the sample standard deviation of the calculated and experimental values. Due to the squaring criterion in the RMSE, larger errors have a greater impact on the MSE than do smaller errors. The mean absolute error (MAE) is the most natural and unambiguous measure of the average error magnitude, reporting the closeness of the prediction to the subsequent results [51]. To describe the average model performance error, we calculate the RMSE, MAE, and normalized RMSE (NRMSE) and



**Fig. 9** The change in the efective threshold voltage of the AlGaN/ GaN HEMT with variation in the pH value, showing a comparison of the data calculated using the present model for devices 2 and 3 versus the experimental data from Ref. [28] and Ref. [24], respectively. The straight and dashed lines show the calculated results, while the solid symbols denote the experimental values

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The RMSE, NRMSE, and MAE values for the  $I_d-V_g$  characteristics of devices 1 and 2 at diferent pH levels are evaluated and summarized in Tables 3 and 4, respectively, revealing that the RMSE of the model proposed herein is 0.053 at pH 4. This is signifcantly better than result obtained using the basic model by Rabbaa et al. [44] with an RMSE value of 0.166 at the same pH level. Similarly, the NRMSE and MAE values of our model verify the closeness of the values calculated herein to the experimental results, as compared with the results in Ref. [44]. The diference between the results calculated using the model in Ref. [44] and the model presented herein is given in Table 4.

Table 5 summarizes the error measurement for the  $I_d-V_d$ characteristics with respect to the experimental results reported in Ref. [24]. The RMSE values presented in Table 5 show the variation due to squaring of larger and smaller errors, while the MAE is constant at 0.001 for the different datasets at different pH levels, thereby demonstrating the





**Table 5** The RMSE values for the  $I_d-V_d$  characteristics of device 3 with solutions at diferent pH levels

$V_{d}$ (V)	Error in $Id$ versus the experimental results in Ref. [24]					
	With solution pH 1.7		With solution pH 11.9			
	At $V_g = 0$ V	At $V_g = -2$ V	At $V_g = 0$ V	At $V_g = -2$ V		
1	0.001	0.00023	0.00011	0.0006		
$\mathcal{D}_{\mathcal{L}}$	0.002	0.001	$-0.0003$	0.001		
3	0.002	0.0005	$-0.0013$	0.001		
4	$-0.0015$	$-0.0015$	$-0.0005$	0.001		
5	$-0.0005$	$-0.003$	0.0006	0.0025		
<b>RMSE</b>	0.0015	0.0015	0.0007	0.0013		
<b>NRMSE</b>	0.061	0.072	0.035	0.095		
MAE	0.001	0.001	0.001	0.001		

reliability of the model presented herein.

## **5 Conclusions**

A new analytical model for the  $I_d - V_d$  characteristics of an AlGaN/GaN HEMT for pH sensing is proposed. The impact of the pH of the electrolyte on the drain current, threshold voltage, and surface potential is examined. The voltage in the gate area is indirectly applied as a function of the ionic concentration. The results demonstrate that, by applying an SiN*x* passivation layer to the HEMT, the sensitivity at lower pH values is improved. The current analytical results show good agreement with experimental results available in literature. The overall average RMSE value of our model is 0.018 for the  $I_d$ – $V_d$  characteristics of different devices, with and without the solution pH. This error range is within

**Table 4** The RMSE values for the  $I_d - V_g$  characteristics of device 2 with solutions at different pH levels

$V_{\rm g}$ (V)	Difference in $I_d$ when calculated using the model in Ref. [44] versus the current model	Error when using the current model to calculate $I_d$ versus the experimental results in Ref. [28]		Error when using the model in Ref. $[44]$ to calculate $I_d$ versus the experimental results in Ref. [28]
	With solution pH 4	With solution pH 4	With solu- tion $pH_1$ 10	With solution pH 4
$-3$	$-0.078$	$-0.072$	0.065	$-0.25$
$-2.5$	$-0.063$	$-0.067$	0.097	$-0.2$
$-2$	$-0.059$	$-0.091$	0.085	$-0.15$
$-1.5$	$-0.086$	$-0.014$	0.007	$-0.1$
$-1$	$-0.084$	0.034	$-0.049$	$-0.05$
$-0.5$		0.003	$-0.056$	$\overline{\phantom{0}}$
$\Omega$		$-0.028$	$-0.008$	
<b>RMSE</b>		0.053	0.061	0.166
<b>NRMSE</b>		0.03	0.037	0.106
MAE		0.044	0.052	0.15

acceptable physical limits. Some limitations were identifed at higher drain voltages, where the calculated and experimental values differ in the  $I_d - V_d$  characteristics. This can be attributed to second-order effects such as short-channel efects (SCE) and self-heating efects (SHE), which are not included in the presented model. SCEs, SHEs, and illumination efects will be included in an extended version of this model in the future.

In summary, the current model exhibits smaller errors and is simpler and more reliable than other models available in the literature. To the best of the authors' knowledge, no such *I*–*V* model has been specifcally implemented for pH sensing applications. In conclusion, the current model will have signifcant applications for device optimization and standardization of prospective chemical and biosensors.

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